

# Introduction to Stokes Structures

IV: explicit computations of one dim'l cases  
via Fourier-Laplace

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## Aim:

- illustrate how the two approaches to the irregular Riemann-Hilbert problem

Stokes-filtered local systems,  
Deligne-Malgrange-Sabbah

Enhanced ind-sheaves,  
d'Agnolo-Kashiwara

cp. Claude's talks

cp. Andrea's talk

can be applied to determine the Stokes structure for some one-dimensional holonomic  $D$ -modules arising as

- the Fourier transform

of a less complicated module

in a topological way (without (multi-)summation)

given as

- linear Stokes data,
- Stokes multipliers

cp. Claude's talk on Monday

## Example of a Fourier transform, applying Deligne-Malgrange-Sabbah's approach

(on a joint work with C. Sabbah, Rend.Sem.Math.Univ.Padova 2015)

Situation:

- $\rho : u \mapsto t = u^p$  a ramification map,
- $\varphi(u) \in u^{-1}\mathbb{C}[u^{-1}]$  an exponential of pole order  $q$ ,
- $R$  a regular singular connection at  $u = 0$  with monodromy data  $(V, T)$ , extended to a free  $\mathbb{C}[u, u^{-1}]$ -module having regular singularity at 0 and  $\infty$
- and

$$\mathcal{M} := \text{El}(\rho, -\varphi, R) := \rho_+(\mathcal{E}^{-\varphi} \otimes R).$$

Consider the Fourier transform  $\widehat{\mathcal{M}}$ , which has a formal structure (Fan, Sabbah)

$$\widehat{\mathcal{M}}^{\wedge\infty} \simeq \text{El}(\widehat{\rho}, \widehat{\varphi}, \widehat{R}),$$

where

- ramification order of  $\widehat{\rho}$  is  $p + q$ ,
- pole order of  $\widehat{\varphi}$  is  $q$  (in particular, have only 1 level assuming  $p, q$  coprime),
- $\widehat{R}$  has monodromy  $(V, (-1)^q T)$ .

Aim:

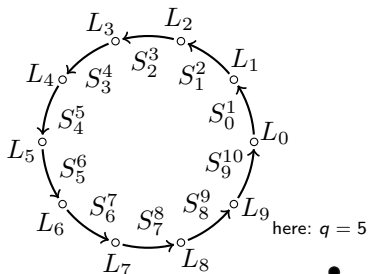
Determine the Stokes structure of  $\widehat{\rho}^+ \widehat{\mathcal{M}}$

(i.e. avoiding ramification. To include it, this would amount to understand  $\mu_{p+q}$ -action)  
in terms of linear Stokes data (recall Claude's talk on Monday):

- vector spaces and isomorphisms as in the picture,
- increasing filtration  $F_\bullet L_{\text{even}}$ ,
- decreasing filtration  $F^\bullet L_{\text{odd}}$ ,
- opposed to each other with respect to the isomorphisms  $S_j^{j+1}$ .

$$L_{2\mu} = \bigoplus_k F_k L_{2\mu} \cap S_{2\mu-1}^{2\mu}(F^k L_{2\mu-1})$$

$$L_{2\mu+1} = \bigoplus_k F^k L_{2\mu+1} \cap S_{2\mu}^{2\mu+1}(F_k L_{2\mu})$$



## Topological computations:

Geometric  $\mathcal{D}$ -module Fourier transform

In the notation

$$\begin{array}{ccc} & \mathbb{A}_t \times \mathbb{G}_{m,\eta} & \\ \pi \swarrow & & \searrow \hat{\pi} \\ \mathbb{A}_t & & \mathbb{G}_{m,\eta} \end{array}$$

we have

$$(\hat{\rho}^+ \widehat{\mathcal{M}})_\infty = \mathbb{C}(\{\eta\}) \otimes_{\mathbb{C}[\eta, \eta^{-1}]} \hat{\pi}_+ (\pi^+ M \otimes E^{-t/\hat{\rho}(\eta)}).$$

Let  $\widehat{\mathcal{L}}$  be the local system of  $\widehat{\mathcal{M}}$  at the circle  $S_\infty^1$  at infinity.

Stokes filtration via moderate deRham complexes

$$\underbrace{\mathrm{DR}^{\mathrm{mod}\infty}(\hat{\rho}^+ \widehat{\mathcal{M}} \otimes \mathcal{E}^{-\hat{\psi}(\eta)})}_{\simeq \widehat{\mathcal{L}}_{\leq \hat{\psi}}}$$

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$$\underbrace{\mathrm{DR}^{\mathrm{mod}\infty}(\widehat{\rho}^+ \widehat{\mathcal{M}} \otimes \mathcal{E}^{-\widehat{\psi}(\eta)})}_{\simeq \widehat{\mathcal{L}}_{\leq \widehat{\psi}}} \xrightarrow{\simeq} R\hat{\pi}_* \mathrm{DR}^{\mathrm{modD}}(\pi^+ \mathcal{M} \otimes E^{-\widehat{\psi}(\eta) - t/\hat{\rho}(\eta)})[1]$$

Isomorphism due to T. Mochizuki.

## Topological computations:

There is a **problem** here:

$$\pi^+ \mathcal{M} \otimes E^{-\widehat{\psi}(\eta) - t/\widehat{\rho}(\eta)}$$

for the necessary choices  $\widehat{\psi}(\eta) = \widehat{\varphi}(\zeta^j \eta)$  for  $\zeta \in \mu_{p+q}$   
(the exponentials of the Fourier transform)

contains exponentials

$$\varphi(u) - \widehat{\psi}(\eta) - \rho(u)/\widehat{\rho}(\eta)$$

with indeterminacies.

Consequently

$$\mathrm{DR}^{\mathrm{modD}}(\pi^+ \mathcal{M} \otimes E^{-\widehat{\psi}(\eta) - t/\widehat{\rho}(\eta)})$$

is not concentrated in one degree and therefore hard to understand (in particular its  $R\tilde{\pi}_*$ ).

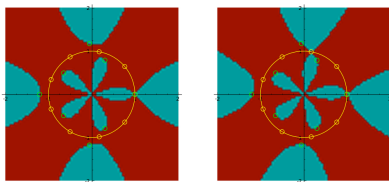
## Blowing-up

these indeterminacies gives rise to a good situation in Mochizuki's sense and hence the possibility to compute the  $R\tilde{\pi}_*$  of a sheaf, one can understand rather easily (and not a complex).

↪ topological computation – even better: can define topological Fourier transform of Stokes-filtered local system compatible with the  $\mathcal{D}$ -module version.

### Example of such a sheaf:

In the case  $p = 4, q = 5$ , typical fibres w.r.t.  $\tilde{\pi}$  are



The sheaf restricted to this fibre, call it  $\mathcal{G}$ , is

- determined by the local system  $\mathcal{L}$  inside the turquoise region and
- zero inside the red region.

$$(\widehat{\mathcal{L}}_{\leq \psi})_{\theta} = H_c^1(\text{fibre over } \theta; \mathcal{G}) \text{ for } \theta \in S^1.$$



## Conclusion for $\mathcal{M} = \text{El}(\rho, \varphi, R)$

- Can compute the linear Stokes data of  $\widehat{\rho}^+ \widehat{\mathcal{M}}$  at  $\infty$  purely topologically (direct images of  $\mathbb{R}$ -constructible sheaves, Leray-covering, ...)
- Example:  $p = 4, q = 5, R = (V, T)$ :

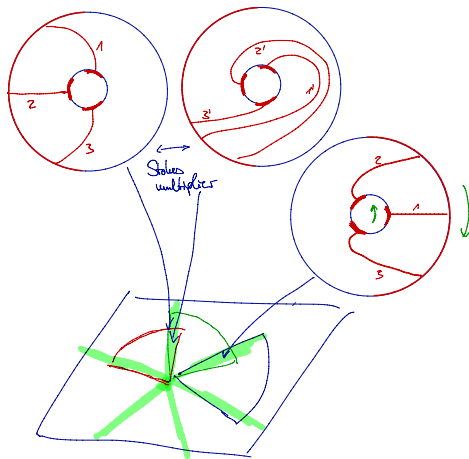
$$L_j := V^{\oplus p+q} =: \bigoplus_{k=0}^{p+q-1} V \otimes \mathbf{1}_k \text{ for all } j,$$

$$F_k L_{2\mu} = \bigoplus_{\nu \leq k} V \otimes \mathbf{1}_\nu$$

$$F^k L_{2\mu+1} = \bigoplus_{\nu \geq k} V \otimes \mathbf{1}_\nu$$



## Another example: Airy



## Example of a Fourier transform, applying d'Agnolo-Kashiwara's approach

d'Agnolo-Kashiwara's R-H correspondence:

Let  $X$  be a complex analytic manifold. Then we have:

$$D_{\text{hol}}^b(\mathcal{D}_X) \xrightarrow{\text{Sol}_X^E} E_{\mathbb{R}\text{-c}}^b(\mathbb{C}_X^{\text{sub}}) \rightarrow D^b(\mathcal{D}_X)$$

fully faithful       reconstruction functor

Important ingredients/constructions:

- bordered spaces
- convolution
- $\mathbb{R}$ -constructibility
- the 'usual' functors like  $Ef_{!!}$ ,  $Ef^{-1}$ , ... and their compatibilities,
- $\mathcal{O}^E$  and hence  $\text{Sol}^E$

## Some definitions/facts:

- $\mathbb{C}_X^E := \varinjlim_{c \rightarrow \infty} \mathbb{C}_{\{t \geq c\}}$

- Fully faithful embedding

$$e : D^b(\mathbb{C}_X^{\text{sub}}) \rightarrow E_{\mathbb{R}-c}^b(\mathbb{C}_X^{\text{sub}}), F \mapsto \mathbb{C}_X^E \otimes \pi^{-1}(F).$$

- $\mathcal{E}_{U/X}^\varphi := \mathcal{D}_X e^\varphi(*Y)$  for  $U \subset X$ ,  $Y = X \setminus U$ .

## Proposition

We have

$$\text{Sol}_X^E(\mathcal{E}_{U/X}^\varphi) \simeq \underbrace{\mathbb{C}_X^E \otimes^+ \mathbb{C}_{\{t + \text{Re } \varphi(x) \geq 0\}}}_{=: E_{U/X}^\varphi} \overset{=: e_{U/X}^\varphi}{\in} E_{\mathbb{R}-c}^b(\mathbb{C}_X^{\text{sub}}).$$

## Fourier transform in various colours

$$\begin{array}{ccccc}
 & \mathbb{A} \times \mathbb{A}^* \times \mathbb{R} & & & \\
 & \swarrow \tilde{p} & \downarrow \bar{p} & \searrow \tilde{q} & \\
 \mathbb{A} \times \mathbb{R} & & \mathbb{A} & & \mathbb{A}^* \times \mathbb{R} \xrightarrow{\tilde{k}} \mathbb{P}^* \times \mathbb{R}
 \end{array}$$

no  $\mathbb{R}$ , no  $\tilde{\phantom{p}}$ ,  $-\phantom{p}$

- $\mathcal{D}$ -module:

$$\begin{aligned}
 \wedge &: D_{\text{hol}}^b(\mathcal{D}_{\mathbb{A}, \infty}) \rightarrow D_{\text{hol}}^b(\mathcal{D}_{\mathbb{A}, \infty}) \\
 \mathcal{M} &\mapsto \widehat{\mathcal{M}} := Dq_*(\mathcal{E}_{\mathbb{A} \times \mathbb{A}^* / \mathbb{P} \times \mathbb{P}^*}^{\text{zw}} \otimes Dp^* \mathcal{M}).
 \end{aligned}$$

- Enhanced sheaves:

$$\begin{aligned}
 \wedge &: \widetilde{E}_{\mathbb{R}\text{-c}}^b(\mathbb{C}_{\mathbb{A} \times \mathbb{R}, \infty}) \rightarrow \widetilde{E}_{\mathbb{R}\text{-c}}^b(\mathbb{C}_{\mathbb{A} \times \mathbb{R}, \infty}) \\
 F &\mapsto F^\wedge := R\tilde{q}_!(e_{\mathbb{A} \times \mathbb{A}^* / \mathbb{P} \times \mathbb{P}^*}^{\text{zw}} \otimes^* \tilde{p}^{-1} F).
 \end{aligned}$$

- Enhanced ind-sheaves:

$$\begin{aligned}
 \wedge &: E_{\mathbb{R}\text{-c}}^b(\mathbb{C}_{\mathbb{A}}^{\text{sub}}) \rightarrow E_{\mathbb{R}\text{-c}}^b(\mathbb{C}_{\mathbb{A}}^{\text{sub}}) \\
 K &\mapsto K^\wedge := Eq_{!!}(E_{\mathbb{A} \times \mathbb{A}^* / \mathbb{P} \times \mathbb{P}^*}^{\text{zw}} \otimes^+ p^{-1} F).
 \end{aligned}$$

# Fourier transform of a perverse sheaf

(on a joint work with A. d'Agnolo, G. Morando and C. Sabbah)

The setting:

- $\mathcal{M}$  a regular singular  $\mathcal{D}$ -module on the affine line  $\mathbb{A}$ , localized at  $\infty$  (notation  $\text{Mod}(\mathcal{D}_{\mathbb{A}, \infty})$ ), with singularities  $\Sigma \subset \mathbb{A}$ ,
- $F := \text{Sol}_{\mathbb{P}}(\mathcal{M})|_{\mathbb{A}}$  the solutions complex, a perverse sheaf on  $\mathbb{A}$ .

Known facts about the Fourier transform  $\widehat{\mathcal{M}}$ :

- regular singular at 0, irregular singular at  $\infty$  and no other singular points,
- the Hukuhara-Levelt-Turritin formal decomposition has the form

$$\widehat{\mathcal{M}}^{\wedge \infty} \simeq \bigoplus_{c \in \Sigma} E_x^c \otimes R_c$$

in the coordinate  $x$  of  $\mathbb{P}^*$  centered at  $\infty$ .

Aim:

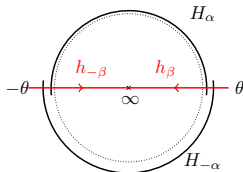
Determine the Stokes structure of the Fourier transform  $\widehat{\mathcal{M}}$  at  $\infty \in \mathbb{P}^*$

Aim:

Determine the Stokes structure of the Fourier transform  $\widehat{\mathcal{M}}$  at  $\infty \in \mathbb{P}^*$ .

We know a priori:

- exponential factors are  $\{\varphi(x) = \frac{c}{x} \mid c \in \Sigma\}$  of slope 1,
- hence, after choosing a starting direction  $\theta \in S_{\infty}^1$  of the real oriented blow-up of  $\mathbb{P}^*$  at  $\infty$ , it suffices to consider two sectors:

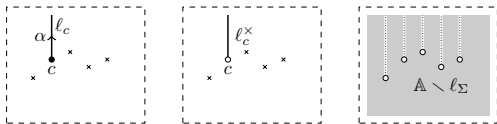




## Linear algebra data associated to $F$ , the quiver of $F$ :

After choice of a suitable pair  $(\alpha, \beta) \in \mathbb{A} \times \mathbb{A}^*$  (fixing a preferred direction/orientation)

- vanishing cycles  $\Phi_c(F) := R\Gamma_c(\mathbb{A}; \mathbb{C}_{\ell_c} \otimes F)$
- (local) nearby cycles  $\Psi_c(F) := R\Gamma_c(\mathbb{A}; \mathbb{C}_{\ell_c^\times} \otimes F)$
- (global) nearby cycles  $\Psi(F) := R\Gamma_c(\mathbb{A}; \mathbb{C}_{\mathbb{A} \setminus \ell_\Sigma} \otimes F)$



and linear maps

$$\Psi(F) \begin{array}{c} \xrightarrow{u_c} \\ \xleftarrow{v_c} \end{array} \Phi_c(F)$$

such that  $1 - uv$  is invertible.

Consider the projections/inclusion

$$\begin{array}{ccccc}
 & & \mathbb{A} \times \mathbb{A}^* \times \mathbb{R} & & \\
 & \swarrow \tilde{p} & \downarrow \bar{p} & \searrow \tilde{q} & \\
 \mathbb{A} \times \mathbb{R} & & \mathbb{A} & & \mathbb{A}^* \times \mathbb{R} \xrightarrow{\tilde{k}} \mathbb{P}^* \times \mathbb{R}
 \end{array}$$

Corollary of to the functorialities in  $E_{\mathbb{R}-c}^b(\mathbb{C}_{\mathbb{P}}^{\text{sub}})$

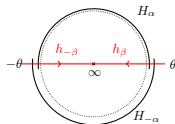
Let  $\mathcal{M} \in D_{rs}^b(\mathcal{D}_{\mathbb{A}, \infty})$  and  $F := \text{Sol}_{\mathbb{P}}(\mathcal{M})|_{\mathbb{A}}$ , then

$$\text{Sol}_{\mathbb{P}^*}^E(\mathcal{M}^\wedge) \simeq \mathbb{C}_{\mathbb{P}^*}^E \otimes^+ \underbrace{R\tilde{k}_! R\tilde{q}!(\mathbb{C}_{\{t+\text{Re}(zw) \geq 0\}} \otimes \bar{p}^{-1}F)[1]}_{\text{complex of usual sheaves on } \mathbb{P}^* \times \mathbb{R}}.$$

Define  $K := R\tilde{q}!(\mathbb{C}_{\{t+\text{Re}(zw) \geq 0\}} \otimes \bar{p}^{-1}F)$ .

## Decomposition in sectors

Let  $H_{\pm\alpha} := \{w \in \mathbb{A}^* \setminus \{0\} \mid \pm \operatorname{Re} \alpha w \geq 0\}$  be the two (closed) sectors and  $H_{\alpha} \cap H_{-\alpha} = h_{\beta} \cup h_{-\beta}$ .



There are natural isomorphisms:

$$s_{\pm\alpha} : \operatorname{Sol}_{\mathbb{P}^*}^E(\widehat{\mathcal{M}})|_{H_{\pm\alpha}} \xrightarrow{\cong} \bigoplus_{c \in \Sigma} (\Phi_c(F) \otimes E^{cW})$$

(where  $\dots|_Y := \pi^{-1}C_Y \otimes \dots$ ).

## Lemma, cp. [d'A-K, last section]

- If  $S$  is a small sector such that  $\operatorname{Re}(cw - dw) > 0$  on  $S$ , then

$$\operatorname{Hom}_{E_{\mathbb{R}-c}^b(\mathbb{C}_{\mathbb{P}^*}^{\text{sub}})}(E^{cw}, E^{dw}) = 0.$$

- If  $S$  contains exactly one Stokes line for each pair  $(c, d)$  in  $\Sigma$  with  $c \neq d$  (e.g.  $S = H_{\pm\alpha}$ ), then

$$\operatorname{End}_{E_{\mathbb{R}-c}^b(\mathbb{C}_{\mathbb{P}^*}^{\text{sub}})}\left(\bigoplus_c \Phi_c \otimes E^{cw}|_S\right) \simeq \mathfrak{t},$$

where

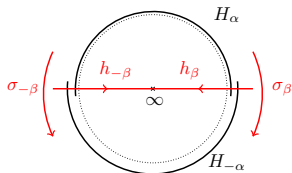
$$\mathfrak{t} := \bigoplus_{c \in \Sigma} \operatorname{End}(\Phi_c) \subset \operatorname{End}\left(\bigoplus_{c \in \Sigma} \Phi_c\right)$$

are the block diagonal matrices.

## Stokes glueing isomorphism

We have both isomorphisms  $s_{\pm\alpha}$  on the common boundary half-lines of  $H_{\pm\alpha}$  and  $K$  is determined by the glueing isomorphisms

$$\sigma_{\pm\beta} := s_{-\alpha}|_{h_{\pm\beta}} \circ (s_{\alpha}|_{h_{\pm\beta}})^{-1} : \bigoplus_{c \in \Sigma} (\Phi_c(F) \otimes E^{cw}) \rightarrow \bigoplus_{c \in \Sigma} (\Phi_c(F) \otimes E^{cw}).$$



Consequences from Lemma 2 slides before:

- the **decomposition isomorphisms**  $s_{\pm\alpha}$  are unique up to base change in  $t$  (block diagonal matrices),
- the **glueing isomorphisms**  $\sigma_{\pm\beta}$  are upper/lower block triangular matrices, notation  $\text{End}^{\pm}$ , with complex coefficients.

Note that  $\alpha$  induces an ordering

$$c_1 <_{\alpha} c_2 <_{\alpha} \dots <_{\alpha} c_n$$

by ordering of  $\text{Re}(\alpha \cdot c)$ .

### Stokes multipliers

We obtain the glueing matrices, the Stokes multipliers

$$S_{\pm\beta} \in \text{End}^{\pm} \left( \bigoplus_{c \in \Sigma} \Phi_c(F) \right).$$

The monodromy of the local system of solutions is

$$T = S_{\beta}^{-1} \cdot S_{-\beta}$$

up to conjugation.

## Topological computation

Recall:

$$\mathrm{Sol}_{\mathbb{P}^*}^E(\mathcal{M}^\wedge) \simeq \mathbb{C}_{\mathbb{P}^*}^E \otimes^+ R\tilde{k}_! \underbrace{R\tilde{q}_!(\mathbb{C}_{\{t+\mathrm{Re}(zw) \geq 0\}} \otimes \bar{p}^{-1}F)}_{=:K}[1].$$

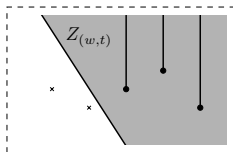
- Consider the stalk of  $K$  at some  $(w, t) \in \mathbb{A}^* \times \mathbb{R}$ :

$$K_{(w,t)} \simeq R\Gamma_c(\mathbb{A}; \mathbb{C}_{Z_{(w,t)}} \otimes F),$$

$$Z_{(w,t)} = \{z \in \mathbb{A} \mid t + \mathrm{Re} zw \geq 0\} = -(t/|w|^2)\bar{w} + \{z \in \mathbb{A} \mid \mathrm{Re} z \geq 0\}\bar{w}$$

a closed half-space,  $\bar{w} := w/|w|$ .

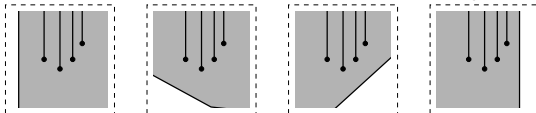
- $Z_{(w,t)} \supset Z_{(w,s)}$  for  $s < t$ .
- $\ell_c \subset Z_{(w,t)} \iff c \in Z_{(w,t)} \iff t + \mathrm{Re} cw \geq 0$ .



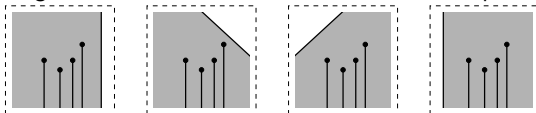
- for  $|w| \gg 0$  and  $t \gg 0$ , we have  $\ell_\Sigma \subset Z_{(w,t)}$ , then

$$K_{(w,t)} \simeq R\Gamma_c(\mathbb{A}; \mathbb{C}_{Z_{(w,t)}} \otimes F) \simeq \bigoplus_{c \in \Sigma} R\Gamma_c(\mathbb{A}; \mathbb{C}_{\ell_c} \otimes F) = \bigoplus_{c \in \Sigma} \Phi_c(F).$$

- Can be globalized to obtain the decomposition isomorphisms  $s_{\pm\alpha}$ .
- The Stokes phenomenon is associated to the following easy observation: rotating  $w$  with  $|w| \gg 0$ , we can use
  - $\ell_c$  as above for the direction  $\alpha$  for  $w$  inside one half-space = sector,



- $\ell_c^-$  using the direction  $-\alpha$  for  $w$  inside the other half-space = sector.





## Result

Result - d'Agnolo, H , Morando, Sabbah

For  $F \in \text{Perv}_\Sigma(\mathbb{C}_A)$  with quiver  $(\Psi, \Phi_i, u_i, v_i)_{c_i \in \Sigma}$ , there is a topological way to compute the Stokes multipliers of its enhanced Fourier-Sato transform and the result is

$$S_\beta = \begin{pmatrix} 1 & u_1 v_2 & u_1 v_3 & \cdots & u_1 v_n \\ & 1 & u_2 v_3 & \cdots & u_2 v_n \\ & & \ddots & & \vdots \\ & & & & 1 \end{pmatrix},$$
$$S_{-\beta} = \begin{pmatrix} \mathbb{T}_1 & & & & \\ -u_2 v_1 & \mathbb{T}_2 & & & \\ -u_3 v_1 & -u_3 v_2 & \ddots & & \\ \vdots & \vdots & & \ddots & \\ -u_n v_1 & -u_n v_2 & \cdots & -u_n v_{n-1} & \mathbb{T}_n \end{pmatrix}.$$

Remark: Cp. the above with

- a general procedure by T. Mochizuki using rapid decay cycles,
- Malgrange's book, chapter XII.

## Monodromy

The monodromy of the Fourier transform  $\widehat{M}$  around  $\infty$  is

$$S_{\beta}^{-1} S_{-\beta} = 1 - \begin{pmatrix} u_1 T_2 T_3 \cdots T_n v_1 & u_1 T_2 T_3 \cdots T_n v_2 & \cdots & u_1 T_2 T_3 \cdots T_n v_n \\ u_2 T_3 \cdots T_n v_1 & u_2 T_3 \cdots T_n v_2 & \cdots & u_2 T_3 \cdots T_n v_n \\ \vdots & \vdots & \ddots & \vdots \\ u_{n-1} T_n v_1 & u_{n-1} T_n v_2 & \cdots & u_{n-1} T_n v_n \\ u_n v_1 & u_n v_2 & \cdots & u_n v_n \end{pmatrix}$$

(This can also be obtained by determining the quiver of  $Sol(\widehat{M})$  at 0).

## Airy function

G.G. Stokes observed the Stokes phenomenon by studying the Airy function, an entire solution to

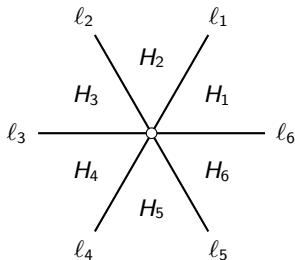
$$(\partial_y^2 - y)u(y) = 0.$$

We have

$$\mathcal{A} := \mathcal{D}_{\mathbb{C}}/\mathcal{D}_{\mathbb{C}}(\partial_y^2 - y) \simeq (\mathcal{E}^{x^3/3})^\wedge$$

Let us study the Stokes structure of the Airy equation as a Fourier transform.

- De-ramify via  $r : \mathbb{C}_v \rightarrow \mathbb{C}_y$ ,  $v \mapsto y = v^2$ , i.e. consider  $r^{-1}\mathcal{A}$ .
- Consider the sectors and their intersections in  $\mathbb{C}_v$  at  $v = \infty$ :



- coordinate change  $\mathbb{C}_u \times \mathbb{C}_v^\times \xrightarrow{\cong} \mathbb{C}_x \times \mathbb{C}_v^\times$  given by

$$\begin{cases} x = \sqrt{-1} uv, \\ v = v, \end{cases} \quad (1)$$

and

$$\begin{aligned} f: \mathbb{C}_u &\rightarrow \mathbb{C}_z = \mathbb{A}, & u &\mapsto u^3 - 3u, \\ g: \mathbb{C}_v &\rightarrow \mathbb{C}_w = \mathbb{A}^*, & v &\mapsto \sqrt{-1} v^3/3, \end{aligned}$$

so that  $xy + x^3/3 = xv^2 + x^3/3 = f(u)g(v) = zw$ .

- Set

$$F := Rf_! \mathbb{C}_{\mathbb{C}_u}[1] \in D^b(\mathbb{C}_{\mathbb{A}}).$$

Recall

$$F := Rf_! \mathbb{C}_{\mathbb{C}_u}[1] \in D^b(\mathbb{C}_A).$$

Then  $F \in \text{Perv}_\Sigma(\mathbb{A})$  for  $\Sigma = \{-2, 2\}$ , and the quiver of  $F$  is

$$\begin{array}{ccc}
 \Phi_2(F) & & \mathbb{C} \\
 \begin{array}{c} \uparrow \\ \nu_2 \\ \downarrow \end{array} & (0 \ 1 \ -1) & \begin{array}{c} \uparrow \\ \downarrow \end{array} \begin{pmatrix} 0 \\ 1 \\ -1 \end{pmatrix} \\
 \Psi(F) \simeq & & \mathbb{C}^3 \\
 \begin{array}{c} \uparrow \\ \nu_{-2} \\ \downarrow \end{array} & (1 \ 0 \ -1) & \begin{array}{c} \uparrow \\ \downarrow \end{array} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} \\
 \Phi_{-2}(F) & & \mathbb{C}.
 \end{array}$$

Key observation

$$Er^{-1}A|_{\mathbb{C}_v^\times} \simeq Eg^{-1}((eF)^\wedge)|_{\mathbb{C}_v^\times}.$$

- Exponential components of  $(eF)^\lambda$  at  $\infty$  are  $E^{\pm 2w}$ ,
- Stokes multipliers are

$$S_\beta = \begin{pmatrix} -1 & 0 \\ -1 & -1 \end{pmatrix}, \quad S_{-\beta} = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}.$$

- $Eg^{-1}E^{\pm 2w} \simeq E^{\pm \frac{2}{3}\sqrt{-1}v^3}$ ,
- $g^{-1}H_\alpha = \bigcup_k H_{2k-1}$ ,  $g^{-1}H_{-\alpha} = \bigcup_k H_{2k}$  and
- hence

$$S_{2k} = S_\beta, \quad S_{2k-1} = S_{-\beta}^{-1}.$$